Tokamak fusion, does it have a path?

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Abstract

stability and steady state plasma regime control, power extraction from both the plasma and the neutron a leap from 1 keV plasma temperature in Russian T-3 machine (1968) to 40 keV and 10.7 MW of DT maintenance of future reactors, etc zone, activation and structural integrity of the machine under 14 MeV fusion neutron bombardment, the way to the tokamak-reactor remain unresolved even at the conceptual level. These problems include the US. Despite much better understanding of the tokamak plasma now, many fundamental problems on fusion power in TFTR at PPPL (1994), are now in an eventual state of defeat and possible shutdown in Tokamak fusion devices, which for 3 decades were leaders in the World fusion program and which made

practical fusion reactor. ing new ideas, which in many aspects rely on the best US tokamak experiments on TFTR (PPPL) and abstraction) leads to breaking with the conventional approach to the tokamak fusion reactor. The resultidea of renewable and absorbing walls at the plasma boundary (which previously was only a theoretical sion for driving liquid metal streams in the tokamak magnetic field. This effect in combination with the DIII-D (GA, San Diego) and basic theory, raise the hope on a new research path for tokamaks toward a This presentation describes the physics of a recently (Dec., 1998) invented method of magnetic propul-



OUTLINE

- 1. Magnetic propulsion of liquid lithium.
- 2. Lithium Walls and tokamak plasma
- (a) LiWalls and energy extraction from the plasma;
- (b) LiWalls and plasma stability;
- (c) LiWalls and plasma energy confinement;
- 3. Yacht Sail approach for tokamak fusion reactors.
- 4. Summary.
- 5. Does tokamak fusion have a path?



Goto Cbpu code.





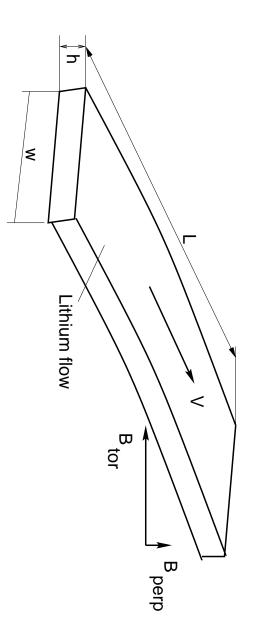
maks There 3 magnetic Reynolds numbers which control lithium MHD in toka-

dynamics: $\Re_0 \equiv \mu_0 \sigma L V$

electro-dynamics: $\Re_1 \equiv \mu_0 \sigma h V,$

 $\mathsf{dynamics}: \;\; \Re_2 \equiv \mu_0 \sigma rac{h^2}{L} V, \ \mu_0 \sigma \simeq 4 \; \left[rac{sec}{m^2}
ight].$

(1.1)



Characteristic flow parameters:

$$V=20~m/sec
ightarrow
ho rac{V^2}{2} \simeq 1~[atm],$$
 $B=1~T
ightarrow rac{B^2}{2\mu_0} = 4~[atm],$ $B=5~T
ightarrow rac{B^2}{2\mu_0} = 100~[atm].$

(1.2)

Dynamic pressure losses are determined by \Re_0 and \Re_2

$$\Re_{0}: \Delta \rho \frac{V^{2}}{2} = \mu_{0} \sigma L V \frac{B_{\perp}^{2}}{2\mu_{0}},$$

$$\Re_{2}: \Delta \rho \frac{V^{2}}{2} = \mu_{0} \sigma \frac{h^{2}}{L} V \Delta \frac{B_{\parallel}^{2}}{2\mu_{0}},$$

$$\mu_{0} \sigma \simeq 4 \left[\frac{sec}{m^{2}} \right].$$
(1)

Magnetic fields from the currents in the stream are determined by \Re_1

$$\Re_1: B_{\parallel out} - B_{\parallel in} = \mu_0 \sigma h V B_{\perp}.$$
 (1.4)

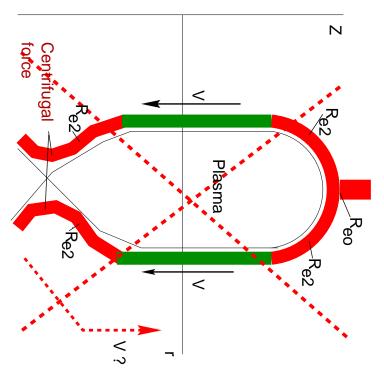


field Lithium "water-falls" will not flow through the tokamak strong toroidal

$$h = 0.1 \ m, \quad L_1 \simeq 0.2 \ m, \quad L_2 \simeq 3 \ m, \quad V > 2 - 5 \ [m/sec],$$

$$\Re_0 = 4L_1V => 1.6,$$

$$\Re_2 = 4\frac{h^2}{L_2}V = 4\frac{h}{L_2}(hV) \simeq 0.01 - 0.025.$$
(1.5)



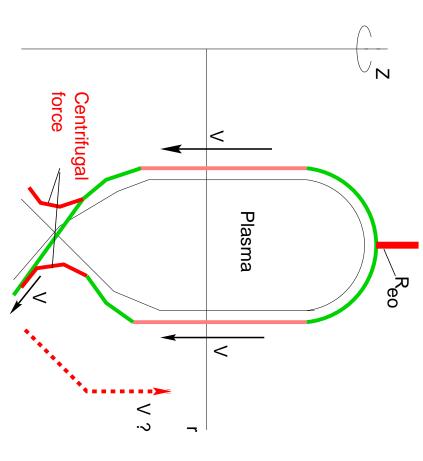




MHD Momentum driven thin walls have many of unresolved problems in lithium

$$h = 0.01 \ m, \quad L_1 \simeq 0.02 \ m, \quad L_2 \simeq 3 \ m, \quad V \simeq 20 \ [m/sec],$$

 $\Re_2 = 4 \frac{h^2}{L_2} V \simeq 1.3 \cdot 10^{-4}.$ (1.6)



$$\Re_0 = 1.6, \quad \rho \frac{V^2}{2} < \Re_0 \frac{B_{pol}^2}{2\mu_0}$$



<u>streams in tokamaks</u> Magnetic propulsion opens the possibility for intense plasma facing lithium

$$p_{\mathbf{j} \times \mathbf{B}}|_{inlet} - p_{\mathbf{j} \times \mathbf{B}}|_{outlet} \gg \Re_2 \frac{B_{tor}^2}{2\mu_0}, \quad \Re_2 \equiv \mu_0 \sigma \frac{h^2}{R} V \simeq 0.0015$$

$$\mu \equiv \mu_0 \sigma \frac{n}{R} V \simeq 0.0015$$

Driving magnetic

pressure

electro-

 $p_{j \times B|outlet} > 1 atm$

 $p_{\mathbf{j} \times \mathbf{B}|inlet} - p_{\mathbf{j} \times \mathbf{B}|outlet} \simeq 1.5$ 3 |atm|

Flow

First wall
Li stream

axis of symmetry

/B_{tor} ,

stream

Pressurized D2

JxB Force

Lithium jets between TF Coils

Toroidal Field Coil

Plasma

Li jets V

Li inlet <u>| 100</u>

Force

Force

He atmosphere

Li stream

parameters

 $V \simeq 20 \ m/sec$ $h \simeq 0.01 \ m$

Magnetic Reynolds numbers

Li jets V

$$\Re_1 \equiv \mu_0 \sigma h V \simeq 0.8, \quad \Re_2 \simeq 0.0015$$

due to centrifugal force Stream are robustly stable

$$ho rac{\langle V^2
angle}{2} > rac{a}{2R} p_{wall} n_r$$

2 Lithium Walls and tokamak plasma

sion reactor between the plasma and the neutron absorbing zone Plasma facing surface is right at the critical place in the fu-

Three things are the most important:

- power and particle extraction;
- effect on core confinement

$$n_{DT} \cdot T_{DT} \cdot \tau_E > 5 \times 10^{21} \ m^{-3} \cdot keV \cdot s, \quad n_{DT} \cdot T_{DT} \cdot \tau_E \propto \tau_E^2$$

3. effect on plasma stability.

$$Power_{fusion} \propto (n_{DT} \cdot T_{DT})^2 \propto \beta^2, \quad \beta \equiv \frac{2\mu_0 \langle 2nT \rangle}{B^2}$$



3 LiWalls and energy extraction from the plasma

tion capabilities Intense lithium streams have reactor relevant power extrac-

$$\Delta T_{max} = q_{wall} \sqrt{rac{4t_{transit}}{\pi \kappa
ho c_p}}, \quad d_{skin} \equiv \sqrt{rac{\kappa t_{transit}}{
ho c_p}}$$

$$R = 6 \ m, \quad a = 1.6 \ m, \quad q_{wall} \simeq 3.5 \ \frac{MW}{m^2}, \quad P_{wall} = 4\pi^2 Raq_{wall} \simeq 1.3 GW$$

even with no reliance on the vortices in the streams.

the neutron wall loading $> 10 MW/m^2$. Intense lithium streams can keep wall temperature low (250-300° C) at



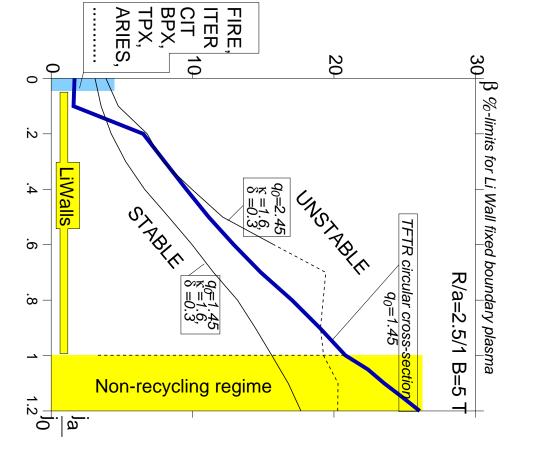
plasma stability control Intense lithium streams may create the best situation for the

$$Power_{fusion} \propto (n_{DT} \cdot T_{DT})^2$$

Goto Cbpu code.



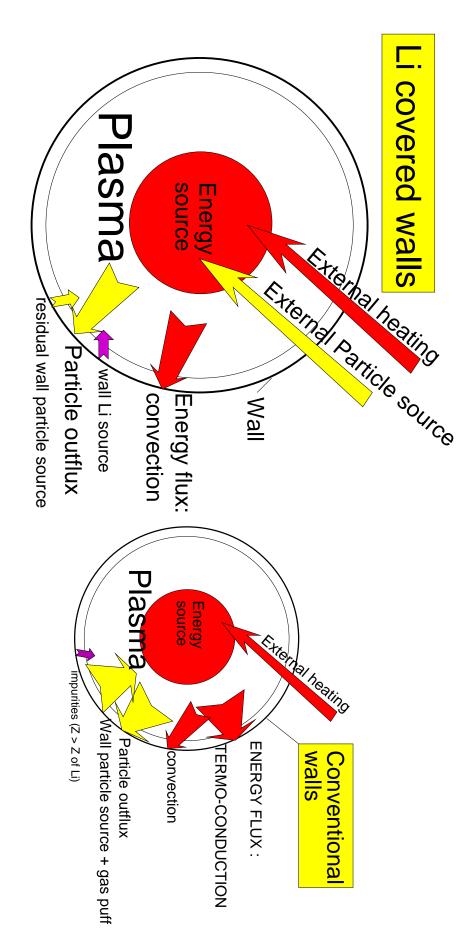
Li walls allow much higher operational plasma pressures.



- eta limits for the second stability regime
- fixed boundary plasma
- n=1,2,3 + ballooning modes (DCON,PEST-2,BALLON)
- current density with an edge pedestal

$$\dot{\mathbf{j}}_{\parallel} = j_a + (j_0 - j_a) \left(1 - rac{r^2}{a^2}\right)$$

Li is an excellent getter for the hydrogen plasma particles.



Lithium can be propelled along the walls for power and particle extraction.



Improved energy confinement is extremely for igniting the plasma

$$n_{DT} \cdot T_{DT} \cdot \tau_E > 5 \times 10^{21} \ m^{-3} \cdot keV \cdot s, \quad n_{DT} \cdot T_{DT} \cdot \tau_E \propto \tau_E^2$$

Plasma profiles are determined by the particle continuity equation

$$\Gamma \equiv Snv = const = (\Gamma)_a$$

and by the energy balance

$$\frac{5}{2}\Gamma T - S(\kappa_T \nabla T + \kappa_n \nabla n) = \int_0^r P_E dv$$

of the (cold) walls and leaves no room for thermo-conduction With perfectly absorbing walls plasma does not know the temperature

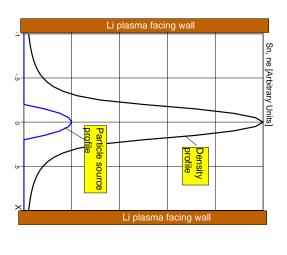
$$\left(rac{5}{2}\Gamma T
ight)_{edge}=\int_0^a P_E dv, \quad T_{edge}=rac{\int_0^a P_E dv}{rac{5}{2}\Gamma} \quad P_E - ext{ heat source}.$$

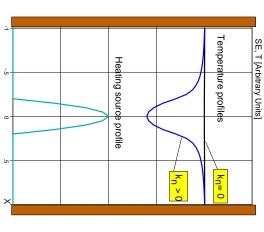
ETON PLASMA CS LABORATOR

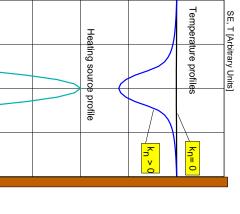
Thus, the major energy loss channel, i.e., thermoconduction, can be eliminated with this absorbing wall boundary condition (S. Krasheninnikov, PFSC at MIT, now at UCSD).





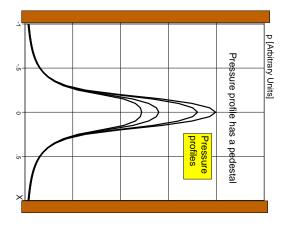


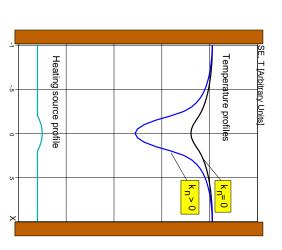




DENSITY profile (left) central fueling. is predetermined by the

the thermo-conduction. order to ELIMINATE (right) adjusts itself in TEMPERATURE profile





plasma boundary. (left) has a jump at the PRESSURE profile

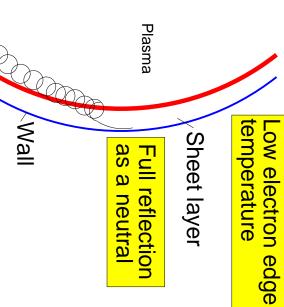
source profile irrespective to the heat thermo-conduction (right) eliminates TEMPERATURE profile

Crucial issues:

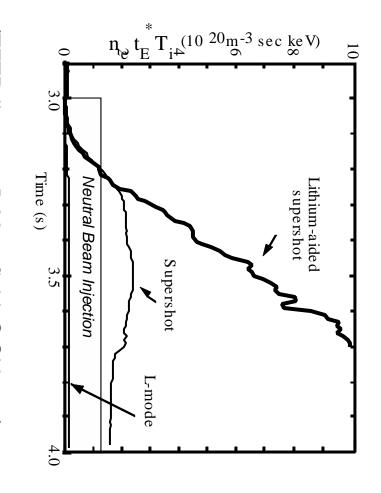
- Properties of lithium. For solid Li they are very attractive, e.g., 1 keV D-ion penetrates hundreds of monlayers of lithium (D.Ruzic, UL).
- Sheet potential near the walls. Is determined by the electron energy, $E \simeq 3T_e/\rho_i$.

Electrons in tokamaks are capable of giving unpleasant surprises.

comparable temperatures are Electron and ion edge Sheet layer by the wall Good absorpsion Wall Plasma



the most important factor in its performance TFTR discovered and demonstrated that Lithium conditioning was F 05



keV 30 keV 30 20 50 10: 20 0 6 40time = TBI + 500 ms # Li pellets # Li pellets 2.8 2.9 3.0 Major Radius (m) Toroidal ITG Modes with Self-Consistent Neoclassical Radial Electric Field SIMULATED $T_{
m i}$ MEASURED $T_{
m i}$ 3.1 3.2 ယ

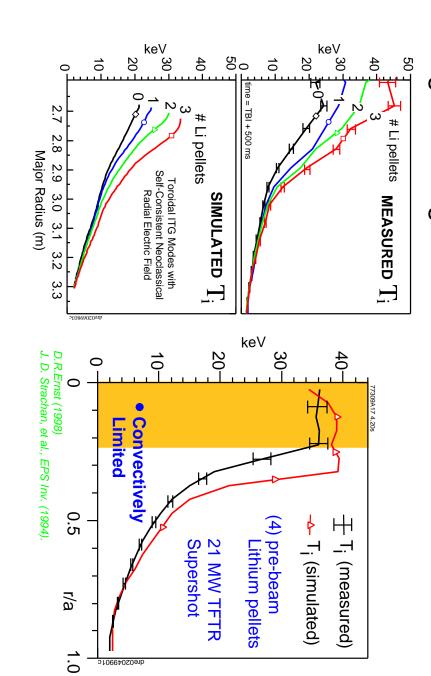
TFTR # 83546 D.Mansfield, C.Skinner)

plasma edge has never been saturated The increase in performance with increase in amount of lithium at the



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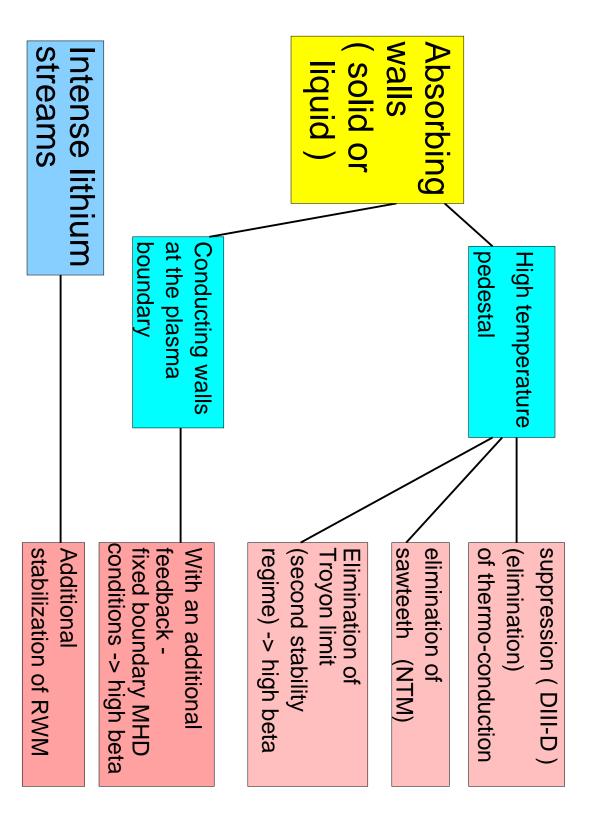
cling at the edge in basic tendencies to flat the temperature in the core by reducing recy-Li Walls concept deffers in details from TFTR results but is consisitent



Goto K.Burrell (DIII-D, GA, CA) talk at APS-2000.



Lithium covered walls affect the very fundamentals of tokamaks

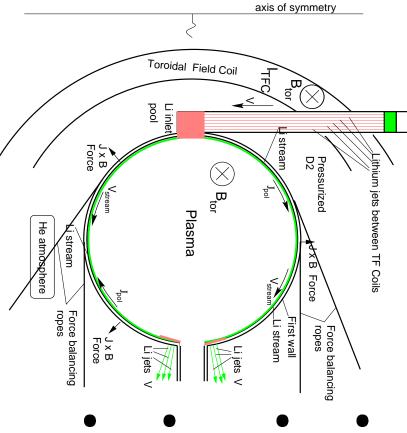




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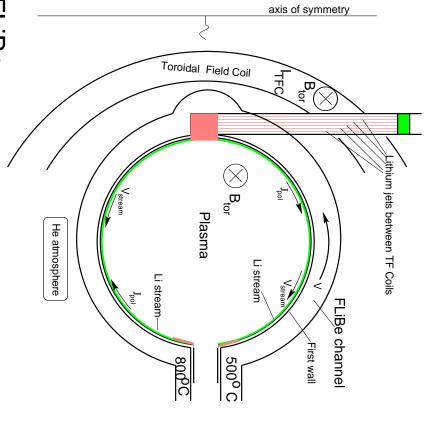
Intense Li Streams affect the very fundamentals of tokamak reactor desing.

Electrodynamic pressure creates a stable situation for the first wall



- Guide wall works against expansion
- Guide wall can be made as a thin shell (like a car tire).
- Toroidal component of the electromagnetic force can be balanced by the set of external wire ropes.
- Inner surface is sealed by the lithium streams (insensitive to cracks) =>>
- Vacuum barrier can be moved to the plasma boundary (giving access to the neutron zone).

combination (S.Zinkle, B.Nelson, ORNL) Intense lithium streams + FLiBe make an excellent FW/blanket



walls.

Lithium streams keep the wall temperature below melting point of FLiBe $\frac{T_{wall} \simeq 200^{\circ} - 250^{\circ} < T_{melt,FLiBe}}{200^{\circ} - 250^{\circ} < T_{melt,FLiBe}}$

Independent of inner temperature in the chanell FLiBe has

15

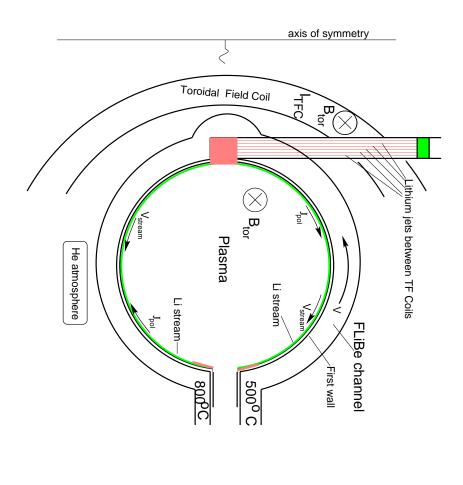
 450^{o}

a solid boundary layer at the

 $(LiF)_n(BeF_2)$

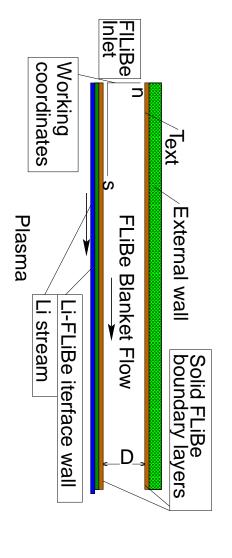
It would be not crazy to think about making the vacuum chamber from the wire mesh





- wall becomes insensitive to drive); thermal deformations (no high-tech for the current pulsed regime is acceptable
- deformations of the wall can be corrected on the fly (Yacht sail approach);
- wire wall presumably can withstand high neutron tlux;
- minimum activation in the neutron zone;
- protection with still excellent coupling plates by the FLiBe layer with the plasma; feedback

Stratified geometry of the FLiBe Blanket/Lithium streams



| 200 | C^{o} | $T_{side\ wall}$ |
|--------|------------------|------------------|
| 100-40 | $\frac{W}{cm^3}$ | S(n) |
| 0.5 | $\frac{m}{sec}$ | V |
| 10 | m | L |
| 0.1 | m | D |

constant by a fast Lithium flow. than the length L of the channel. Plasma side wall temperature is kept The radial thickness D of the channel is assumed to be much smaller

Heat source S corresponds approximately to 10 MW/m 2 in neutrons.



two solid salt layers are formed on the walls of the channel. The stationary heat diffusion equation The walls of the channel are kept below the melting point of FLiBe, so

$$\rho c_p V \frac{\partial T}{\partial s} = \kappa T_{nn}'' + S, \quad T > T_{melt},$$

$$0 = \kappa T_{nn}'' + S, \quad T < T_{melt}$$
(5.1)

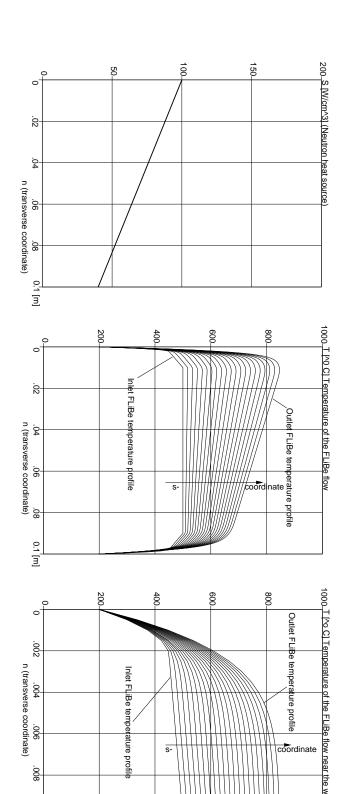
together with the matching conditions determines the temperature distribution in the flow

velocity of the flow, κ is the thermo-conduction. Here, ho is the mass density of FLiBe, c_p is the heat capacity, V is the

lem in a self-consistent way. Thickness of the solid layer is determined as an eigenvalue of the prob-

| $rac{kg}{m^3}$ 2240 $rac{g \cdot Co}{q \cdot Co}$ 2380 $rac{G}{n \cdot Co}$ 450 |
|--|
|--|

Profiles of the (neutron) heat source and T in the FLiBe channel



of the flow is determined solely by the heat source power FLiBe thermo-conduction is so small that the temperature inside body

$$\rho c_p V \frac{\partial T}{\partial s} \simeq S, \quad T > T_{melt},$$
(5.2)

not by thermo-conduction losses.



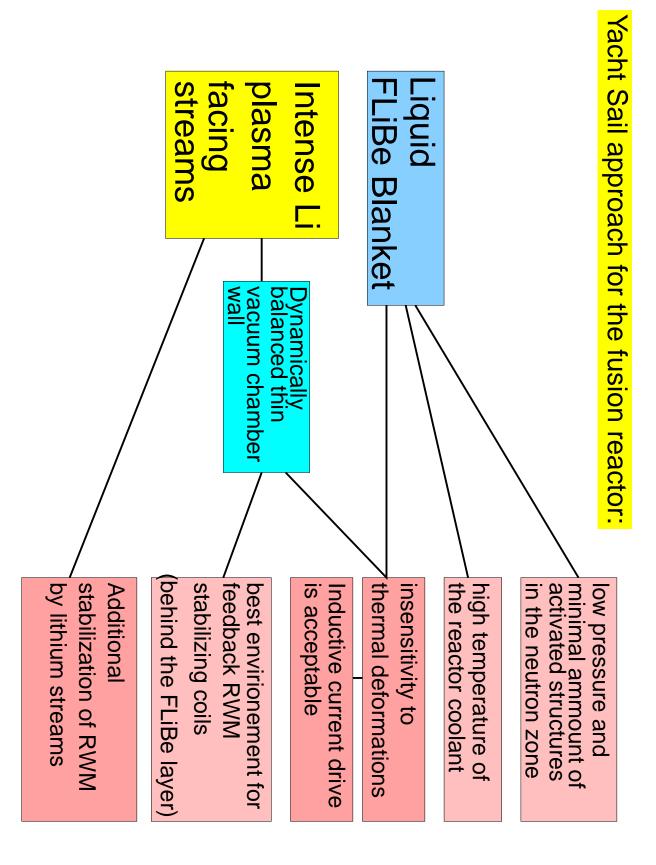
the channel. Inside, each of them contains a sublayer of a solid FLiBe. Two boundary layers of the order of 1-3 mm are formed near walls of

side of the wall, which constitute approximately 4 % of the incoming neutron flux energy. plasma side wall and $0.16\ MW/m^2$ through the Toroidal Field Coil (TFC) In the example the averaged energy losses are $0.26\ MW/m^2$ through the

FLiBe seems to be a perfect coolant for the tokamak-reactor







6 Summary

as a consistent tokamak-rector concept. 1999), the idea of a tokamak with lithium covered walls was presented From the very first presentation to PPPL (PS&T seminar, Jan.

introduced into the tokamak research For the first time, the renewable and absorbing plasma facing walls were

From the plasma physics side, lithium walls may provide

- a low recycling regime (best possible for energy confinement);
- low-Z plasma facing surface (with a central plasma fueling and face impurities source);
- rector relevant power extraction capabilities from the plasma
- wall conditions, which are not sensitive to the edge plasma temperature as soon as it exceeds a certain level (about 1 keV).
- slowing down free-boundary MHD instabilities,
- etc,



fusion reactor design, which may provide Li Walls, for the first time, introduce the "Yacht Sail" approach for the

- insensitivity of the structure to thermal/electro-magnetic perturba-
- best environment for both internal and external plasma stability con-
- elimination of (unrealistic) requirement for the stationary regime;
- efficient power extraction from the neutron zone with a high temperature (FLiBe) coolant;
- minimizing the content of activated structural elements;
- simplification of the entire reactor control and maintainence scheme
- . . .



DETON PLASMA SICS LABORATORY Leonid E. Zakharov, PPPL Theory seminar, PU, Jan.18, 2001, Princeton NJ

7 Does the tokamak fusion have a path?

Go to netscape



meaning of what Sean Connery said about the tokamak fusion Inside the fusion, we, fusion physicists, should, first, realize the deep

"It is impossible (in the conventional way, LZ)"

... but doable (if we follow the way of physics, LZ)"

REGENCY ENTERPRISES film "Entrapment", (S. Connery, TWENTIETH CENTURY FOX and 1999)

be on the right track in the great endeavor to make the tokamak working in our life time limit as a commercial power reactor. You, our physics colleagues outside fusion, should always remind us to



